

AMATH 402
Introduction to Dynamical Systems and Chaos
Winter 2009
Homework 3: Solutions

Due: Wednesday, January 28, 2009

1. 4.1.1

$\sin(a\theta)$ has to be 2π periodic. This means that $a \in \mathbb{Z}$, i.e., an integer.

2. 4.1.2 Equation

$$\theta' = 1 + 2 \cos \theta$$

has fixed points at

$$\cos \theta^* = -\frac{1}{2} \Rightarrow \theta_1^* = \frac{2\pi}{3} \text{ stable, } \theta_2^* = \frac{4\pi}{3} \text{ unstable.}$$



3. 4.1.4 Equation

$$\theta' = \sin^3 \theta$$

has fixed points at

$$\sin^3 \theta^* = 0.$$

That is $\theta_1^* = 0$ which is unstable, and $\theta_2^* = \pi$ which is stable.

4. 4.1.6 Equation

$$\theta' = 3 + \cos(2\theta)$$

has its right-hand-side

$$3 + \cos(2\theta) > 0 \quad \forall \theta \in [0, 2\pi).$$

Therefore, there is no fixed point on the cycle. The motion according to the equation is a continuous rotation.

5. 4.3.2

(a)

$$\begin{aligned} u &= \tan\left(\frac{\theta}{2}\right), \\ \theta &= 2 \arctan u, \\ d\theta &= \frac{2du}{1+u^2}. \end{aligned}$$

One can also obtain the last equation by

$$\begin{aligned} du &= d\left\{\tan\left(\frac{\theta}{2}\right)\right\} = \sec^2\left(\frac{\theta}{2}\right) \frac{d\theta}{2}, \\ d\theta &= 2 \sec^{-2}\left(\frac{\theta}{2}\right) du = \frac{2du}{(1+\tan^2(\theta/2))} = \frac{2du}{(1+u^2)} \end{aligned}$$

(b)

$$\sin \theta = 2 \sin \frac{\theta}{2} \cos \frac{\theta}{2} = 2 \tan \frac{\theta}{2} \cos^2 \frac{\theta}{2} = \frac{2 \tan(\theta/2)}{1 + \tan^2(\theta/2)} = \frac{2u}{1+u^2}.$$

(c)

$$\theta = \pm\pi \Rightarrow \frac{\theta}{2} = \pm\frac{\pi}{2} \Rightarrow u = \tan \frac{\theta}{2} = \tan\left(\pm\frac{\pi}{2}\right) = \pm\infty.$$

Therefore, the integration

$$\int_{-\pi}^{\pi} \frac{d\theta}{\omega - a \sin \theta}$$

can be written as

$$\int_{-\infty}^{\infty} \frac{\frac{2du}{1+u^2}}{\omega - a \frac{2u}{1+u^2}} = \int_{-\infty}^{\infty} \frac{2du}{\omega(1+u^2) - 2au}.$$

(d),(e) Therefore, the T on page 98

$$\begin{aligned} T &= \int_{-\infty}^{\infty} \frac{2du}{\omega(1+u^2) - 2au} \\ &= \frac{1}{\omega} \int_{-\infty}^{\infty} \frac{2du}{1+u^2 - 2au/\omega} \\ &= \frac{1}{\omega} \int_{-\infty}^{\infty} \frac{2du}{u^2 - 2au/\omega + (a/\omega)^2 + 1 - (a/\omega)^2} \\ &= \frac{2}{\omega} \int_{-\infty}^{\infty} \frac{d(u - a/\omega)}{(u - a/\omega)^2 + (1 - (a/\omega)^2)} \end{aligned}$$

According to the result in 4.3.1, if $a/\omega < 1$, we thus have

$$T = \frac{2\pi}{\omega\sqrt{1 - (a/\omega)^2}} = \frac{2\pi}{\sqrt{\omega^2 - a^2}}.$$

6. 4.3.3

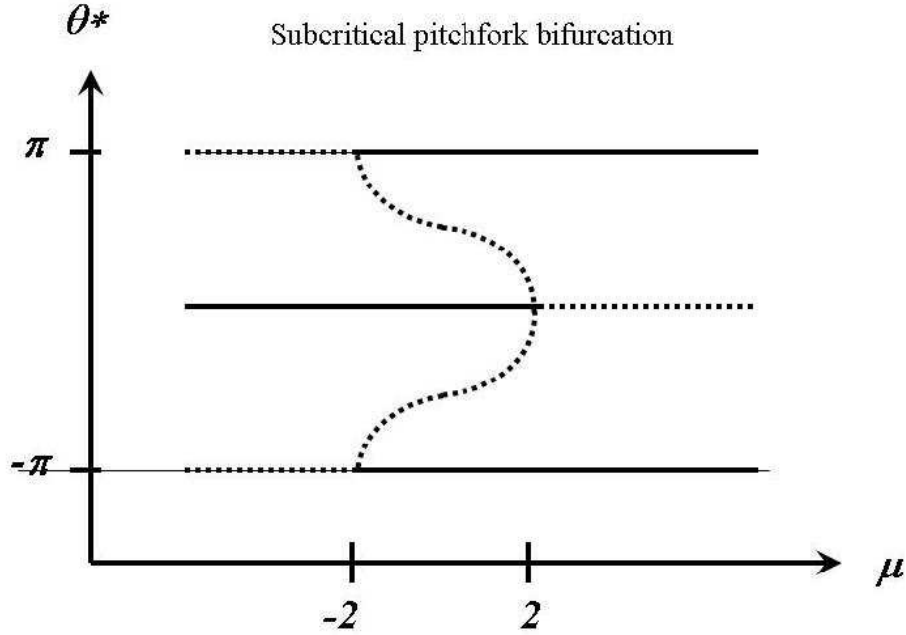
Equation

$$\begin{aligned} \theta &= \mu \sin \theta - \sin(2\theta) \\ &= \mu \sin \theta - 2 \sin \theta \cos \theta \\ &= \sin \theta (\mu - 2 \cos \theta). \end{aligned}$$

The fixed points are at

$$\begin{aligned} \sin \theta = 0, &\Rightarrow \theta = 0, \pi; \\ \cos \theta = \frac{\mu}{2}, &\Rightarrow \theta = \pm \arccos\left(\frac{\mu}{2}\right) \text{ if } |\mu| < 2 \end{aligned}$$

At $\mu = \pm 2$ there are bifurcation values.



7. 4.3.4 Equation

$$\theta' = \frac{\sin \theta}{\mu + \cos \theta}$$

has both fixed point(s) where $\theta' = 0$ and singularity where $\theta' = \infty$.

For $|\mu| > 1$, the denominator $\mu + \cos \theta$ does not change sign. Hence, the fixed points are at $\sin \theta = 0$. They are $\theta_1^* = 0$ and $\theta_2^* = \pi$.

$$\theta_1^* = 0 \text{ stable, and } \theta_2^* = \pi \text{ unstable if } \mu < -1,$$

$$\theta_1^* = 0 \text{ unstable, and } \theta_2^* = \pi \text{ stable if } \mu > 1.$$

For $|\mu| < 1$, there are two roots for the denominator $\mu + \cos \theta = 0$, $\Rightarrow \theta_1^\ddagger$ and θ_2^\ddagger . We shall denote $\theta_1^\ddagger \in (0, \pi)$ and $\theta_2^\ddagger \in (\pi, 2\pi)$.

We now analyze the singularity θ_1^\ddagger . For $\theta < \theta_1^\ddagger$, $\frac{\sin \theta}{\mu + \cos \theta} > 0$ because the numerator is positive and $\cos \theta$ is a decreasing function. But for $\theta > \theta_1^\ddagger$, $\frac{\sin \theta}{\mu + \cos \theta} < 0$. Therefore, the singularity θ_1^\ddagger is an attractor. Near the singularity θ_2^\ddagger , $\frac{\sin \theta}{\mu + \cos \theta} > 0$ for $\theta < \theta_2^\ddagger$ and $\frac{\sin \theta}{\mu + \cos \theta} < 0$ for $\theta > \theta_2^\ddagger$, so again the singularity θ_2^\ddagger is also an attractor (see figure in which the black squares represent the singularities.)

For $\mu = -1$, the fixed point $\theta_2^* = \pi$ is unstable. The fixed point $\theta_1^* = 0$ and singularities θ_1^\ddagger , θ_2^\ddagger merge:

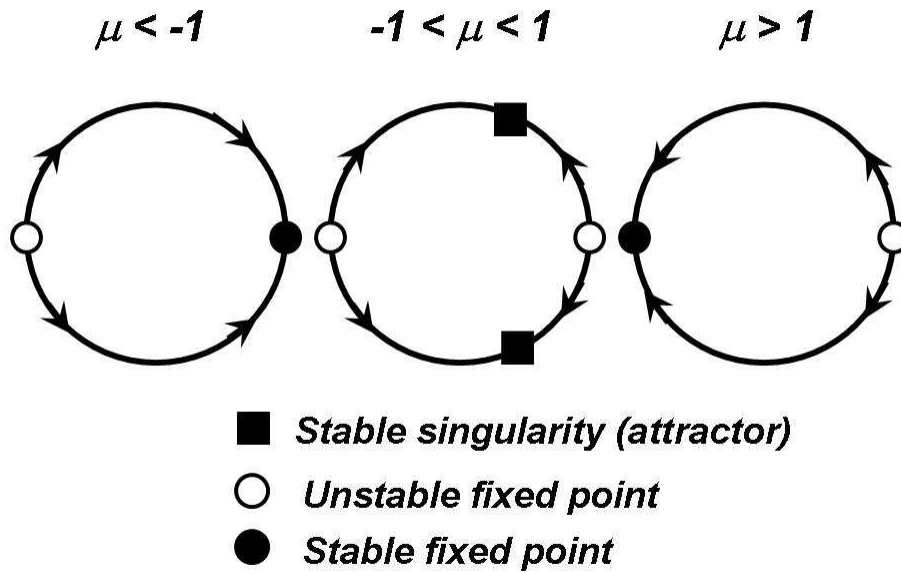
$$\frac{\sin \theta}{-1 + \cos \theta} \approx \frac{\theta - \theta^3/3!}{-\theta^2/2 + \dots} \approx -\frac{2}{\theta},$$

which is an attractive singular point.

For $\mu = +1$, the fixed point $\theta_1^* = 0$ is unstable. fixed point $\theta_2^* = \pi$ and singularity $\theta_1^\ddagger, \theta_2^\ddagger$ merge. Let $\phi = \theta - \pi$, then

$$\frac{\sin \theta}{1 + \cos \theta} = \frac{\sin(\phi + \pi)}{1 + \cos(\phi + \pi)} = \frac{-\sin(\phi)}{1 - \cos(\phi)} \approx \frac{-\phi + \phi^3/3!}{\phi^2/2 + \dots} \approx -\frac{2}{\phi},$$

which is again an attractive singular point.



8. 4.3.5 Equation

$$\theta' = \mu + \cos \theta + \cos(2\theta)$$

can be re-written as

$$\theta' = (\mu - 1) + \cos \theta + 2 \cos^2 \theta. \quad (1)$$

If $\mu > 9/8$, i.e., $\mu - 1 > 1/8$, then $f(x) = \mu - 1 + x + 2x^2 > 0$ for all x . Hence the right-hand-side of Eq. (1) is positive for all θ . There is no fixed point; the motion is a continuous counter-clockwise rotation.

If $\mu < -2$, i.e., $\mu - 1 < -3$, then the two roots of $f(x)$, x_1^* and $x_2^* \notin [-\frac{3}{2}, 1] \Rightarrow$ no fixed point for Eq. (1). The motion is a continuous clockwise rotation.

For $-2 < \mu < 0$, the larger one of the two roots of $f(x)$, $x_2^* \in (\frac{1}{2}, 1)$, while the $x_1^* < -1$:

$$x_2^* = \frac{-1 + \sqrt{1 - 8(\mu - 1)}}{4}.$$

Therefore, Eq. (1) has two fixed points

$$\cos \theta = x_2^*.$$

For $0 < \mu < 9/8$, both roots of $f(x)$, x_1^*, x_2^* are now $\in (-1, 1)$: $-1 < x_1^* < -\frac{1}{4} < x_2^* < 1$. This corresponds to four fixed points.

At $\mu = -2$, there is a saddle-node bifurcation at $\theta = 0$ (2π). **At** $\mu = 0$, there is a saddle-node bifurcation at $\theta = \pi$. **At** $\mu = 5/8$, there are two simultaneous saddle-node bifurcations, at $\arccos(-\frac{1}{4})$, they are at $\theta = 104.5^\circ$ and $\theta = 225.5^\circ$.

