

AMath 403 Spring 2004
Final exam review

1 Course review

1. Separation of variables; Fourier series, Fourier sine series, Fourier cosine series; eigenfunction expansion

- (a) These methods are used on a *finite domain*.
- (b) The solution is separated into functions of single variables, e.g., $u(x, t) = X(x)T(t)$. The PDE is then written as multiple ODEs. (This is not always possible!) The ODE with homogeneous boundary conditions is called the *eigenvalue problem*. (This is usually in x .) The general solution is written then as a sum of eigenfunctions with coefficients chosen to satisfy the non-homogeneous boundary or initial conditions.
- (c) For homogeneous Dirichlet boundary conditions, i.e., $u(0, t) = u(L, t) = 0$, the eigenfunction are sines. This gives *Fourier sine series*. The Fourier sine series for $f(x)$ on $0 < x < L$ is

$$f(x) = \sum_{n=1}^{+\infty} b_n \sin\left(\frac{n\pi x}{L}\right),$$
$$b_n = \frac{2}{L} \int_0^L f(x) \sin\left(\frac{n\pi x}{L}\right) dx.$$

- (d) For homogeneous Neumann boundary conditions, i.e., $\frac{\partial u}{\partial x}(0, t) = \frac{\partial u}{\partial x}(L, t) = 0$, the eigenfunction are cosines. This gives *Fourier cosine series*. The Fourier cosine series for $f(x)$ on $0 < x < L$ is

$$f(x) = \sum_{n=0}^{+\infty} a_n \cos\left(\frac{n\pi x}{L}\right),$$
$$a_0 = \frac{1}{L} \int_0^L f(x) dx,$$
$$a_n = \frac{2}{L} \int_0^L f(x) \cos\left(\frac{n\pi x}{L}\right) dx, \quad \text{for } n = 1, 2, \dots$$

- (e) For periodic boundary conditions, i.e., $u(-L, t) = u(L, t)$ and $\frac{\partial u}{\partial x}(-L, t) = \frac{\partial u}{\partial x}(L, t)$, the eigenfunctions are sines and cosines. This gives *Fourier series*. The Fourier series for $f(x)$ on $-L < x < L$ is

$$f(x) = \sum_{n=0}^{+\infty} a_n \cos\left(\frac{n\pi x}{L}\right) + b_n \sin\left(\frac{n\pi x}{L}\right),$$

$$a_0 = \frac{1}{2L} \int_{-L}^L f(x) \, dx,$$

$$a_n = \frac{1}{L} \int_{-L}^L f(x) \cos\left(\frac{n\pi x}{L}\right) \, dx, \quad \text{for } n = 1, 2, \dots,$$

$$b_n = \frac{1}{L} \int_{-L}^L f(x) \sin\left(\frac{n\pi x}{L}\right) \, dx.$$

- (f) For *non-homogeneous* PDEs, use eigenfunction expansion. That is, write every function $f(x, t)$ in a series of the eigenfunctions for the *homogeneous* problem, ϕ_n , i.e., $f(x, t) = \sum \alpha_n(t) \phi_n(x)$. This gives an ODE for $\alpha_n(t)$
- (g) All of these methods can only deal with *one non-homogeneity* at a time! To solve a problem with multiple non-homogeneity, solve the problems with only *one non-homogeneity* at a time and then add the solutions.

2. Higher-dimensional PDE; Sturm–Liouville problems

- (a) Use the same idea as for two-dimensional problems but several times, e.g., $u(x, y, t) = U(x, y)T(t)$ and then $U(x, y) = X(x)Y(y)$. This gives n ODEs for an n -dimensional PDE.
- (b) For *non-rectangular* coordinates the eigenfunctions are not sines and cosines. Sturm–Liouville theory says solutions to the PDE can be written as sums of the eigenfunctions.
- (c) *Polar coordinates* and *cylindrical coordinates* give Bessel’s equation,

$$\frac{d}{dr} \left(r \frac{dR}{dr} \right) + \left(\lambda r - \frac{m^2}{r} \right) R = 0.$$

The solutions are Bessel functions, $J_m(\sqrt{\lambda}r)$ and $Y_m(\sqrt{\lambda}r)$. (Y_m is unbounded at $r = 0$ so only J_m is used on problems where $r = 0$ is inside the domain.)

- (d) *Spherical coordinates* give the spherical Bessel's equation,

$$\frac{d}{d\rho} \left(\rho^2 \frac{dR}{d\rho} \right) + [\lambda\rho^2 - m(m+1)] R = 0.$$

The solutions are spherical Bessel functions, $j_m(\sqrt{\lambda}r)$ and $n_m(\sqrt{\lambda}r)$. (n_m is unbounded at $\rho = 0$ so only j_m is used on problems where $r = 0$ is inside the domain; we didn't talk about n_m in class, but it analogous to Y_m for polar coordinates.)

3. Fourier transforms

- (a) This method is used on an *infinite domain*; it is the analog of Fourier series.
- (b) The Fourier transform and inverse Fourier transform are defined by

$$\mathcal{F}[f] = \hat{f}(\omega) = \int_{-\infty}^{+\infty} f(x) e^{i\omega x} dx,$$

$$\mathcal{F}^{-1}[\hat{f}] = f(x) = \frac{1}{2\pi} \int_{-\infty}^{+\infty} \hat{f}(\omega) e^{-i\omega x} d\omega.$$

- (c) The Fourier transform of derivatives gives

$$\mathcal{F} \left[\frac{\partial^n u}{\partial x^n} \right] = (-i\omega)^n \hat{u},$$

$$\mathcal{F} \left[\frac{\partial^n u}{\partial t^n} \right] = \frac{\partial^n \hat{u}}{\partial t^n}.$$

- (d) The Fourier transform of a PDE gives an ODE. Solving the ODE and then computing the inverse Fourier transform gives the solution.

4. Green's functions

- (a) Green's functions can be used on either finite or infinite domains. Green's functions give the solution to a PDE in terms of the integral of non-homogeneities with the Green's function.
- (b) To problem to solve to find the Green's function is the *adjoint problem*, with both the *adjoint operator* and the *adjoint* boundary and initial conditions (which are found from the *homogeneous* boundary and initial conditions on the original problem).
- (c) On infinite domains, Green's functions can be found by radial symmetry or Fourier transforms. These are called *free-space Green's functions*. Often we just look these up in books.
- (d) The *method of images* finds Green's functions for *semi-infinite* or *finite* domains from free-space Green's functions by geometric arguments.
- (e) Eigenfunction expansions can be used to find Green's functions on *finite* domains.
- (f) Boundary and initial conditions can be incorporated by looking at the remaining terms from computing the adjoint.

5. Characteristics

- (a) The method of characteristics is used for solving *first-order quasi-linear* PDE, i.e.

$$\frac{\partial u}{\partial t} + c(u, x, t) \frac{\partial u}{\partial x} = f(u, x, t).$$

- (b) The method is to look for curves $(x(t), t)$ on which the PDE becomes an ODE. This results in the system of ODE

$$\begin{aligned} \frac{du}{dt} &= f(u, x, t), \\ \frac{dx}{dt} &= c(u, x, t), \end{aligned}$$

with the initial conditions $u(x_0, 0) = u_0(x_0)$ and $x(0) = x_0$.

- (c) The solution for x , $x = X(t, x_0)$, gives the *characteristic curves*. If the curves do not intersect, the solution is valid for all time. The point at which the curves intersect can be found by look for the first point at which $\frac{\partial u}{\partial x}$ or $\frac{\partial u}{\partial t}$ is infinite.

- (d) If the characteristic curves intersect, a *shock solution* can be constructed using an *integral conservation law*.
- (e) If there are no characteristics in a region, a solution can be constructed by inserting a *fan* into the region.

2 Exam review problems

1. Solve the Schrödinger equation in 1 spatial dimension,

$$i\hbar \frac{\partial u}{\partial t} = -\frac{\hbar^2}{2\mu} \frac{\partial^2 u}{\partial x^2}, \quad 0 < x < L, \quad t > 0,$$

with

$$u(0, t) = u(L, t) = 0,$$

and an arbitrary initial condition.

2. Solve the Schrödinger equation in 2 spatial dimensions,

$$i\hbar \frac{\partial u}{\partial t} = -\frac{\hbar^2}{2\mu} \left[\frac{1}{r} \frac{\partial}{\partial r} \left(r \frac{\partial u}{\partial r} \right) + \frac{1}{r^2} \frac{\partial^2 u}{\partial \theta^2} \right], \quad 0 < r < a, \quad t > 0,$$

with

$$u(a, \theta, t) = 0,$$

and an arbitrary initial condition.

3. Solve

$$\frac{\partial u}{\partial t} = \alpha^2 \frac{\partial^2 u}{\partial x^2}, \quad -\infty < x < +\infty, \quad t > 0,$$

with

$$u(x, 0) = \delta(x).$$

4. Find the Green's function solution for

$$\frac{\partial u}{\partial t} - \frac{\partial^2 u}{\partial x^2} = f(x, t), \quad 0 < x < +\infty, \quad t > 0,$$

with

$$u(0, t) = g(t),$$

and

$$u(x, 0) = h(x),$$

given that the free-space Green's function is

$$G_{FS}(x, t; x_0, t_0) = \frac{H(t_0 - t)}{\sqrt{4\pi(t_0 - t)}} e^{-\frac{(x-x_0)^2}{4(t_0-t)}}.$$

First, find the Green's function for the semi-infinite domain by the method of images and then figure out what the extra terms are needed for the non-homogeneous boundary and initial conditions.

5. Solve

$$\frac{\partial u}{\partial t} + u^2 \frac{\partial u}{\partial x} = \frac{1}{u}, \quad -\infty < x < +\infty, \quad t > 0,$$

with

$$u(x, 0) = 0.$$